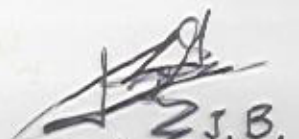


NNN'02

PHENOMENOLOGY of ν ' OSCILLATIONS

- Consisting Evidence for
 ν ' Masses & Mixings
from Atmospheric & Solar.
- Established Facts if
LBL, Reactor confirm.
- ν_{e3} ?, ~~CP~~?
- Sign [Δm^2], Absolute Mass
- Dirac or Majorana?


J.B.
U. Valencia.

Neutrinos as a Unique Probe: $10^{-33} - 10^{+28}$ cm

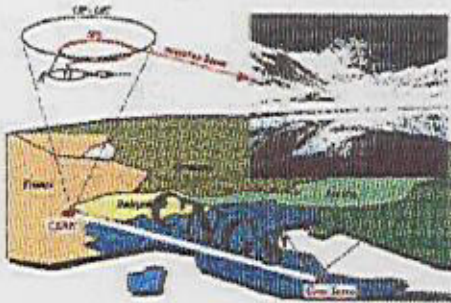
● Particle Physics

- $\nu N, \mu N, eN$ scattering: existence/ properties of quarks, QCD
- Weak decays ($n \rightarrow pe^- \bar{\nu}_e, \mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$): Fermi theory, parity violation, mixing
- Neutral current, Z-pole, atomic parity: electroweak unification, field theory, m_t ; severe constraint on physics to TeV scale
- Neutrino mass: constraint on TeV physics, grand unification, superstrings

● Astrophysics/Cosmology

- Core of Sun
- Supernova dynamics
- Atmospheric neutrinos (cosmic rays)
- AGNs, cosmic rays, violent events, GRB
- Large scale structure (dark matter)
- Nucleosynthesis (big bang - small A ; stellar - to iron; supernova - large N)
- Baryogenesis
- Simultaneous probes of ν and astrophysics

CERN to Gran Sasso Neutrino Beam



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- NEUTRINO OSCILLATIONS

For 3 families of active massive ν 's, with mixing,

$$\nu_\alpha = \sum_i U_{\alpha i} \nu_i + \dots$$

[ν_i are either D or M] \leftarrow sterile later

- Flavour Evolution in Vacuum

$$|\nu_\alpha(t)\rangle = \sum_\beta |\nu_\beta\rangle \left(\sum_i U_{\beta i} e^{-iE_i t} U_{\alpha i}^* \right) + \dots$$

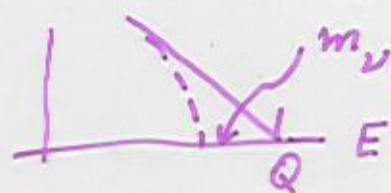
- Transition Probability

$$P(\nu_\alpha \rightarrow \nu_\beta) = \delta_{\alpha\beta} + 2 \operatorname{Re} \sum_{j < k} U_{\beta j} U_{\alpha j}^* U_{\beta k}^* U_{\alpha k} \left(e^{-i \Delta m_{jk}^2 \frac{L}{p}} - 1 \right)$$

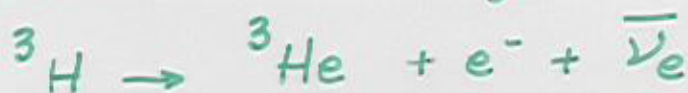
- $\Delta m_{jk}^2 \frac{L}{p} \ll 1, \forall (j, k) \Rightarrow P(\nu_\alpha \rightarrow \nu_\beta) \rightarrow \delta_{\alpha\beta}$

DIRECT MEASUREMENT OF ν -MASS

Reminder: If $m_\nu \neq 0$, the hard part of β -spectra will be distorted \rightarrow plot



The "classical" decay is



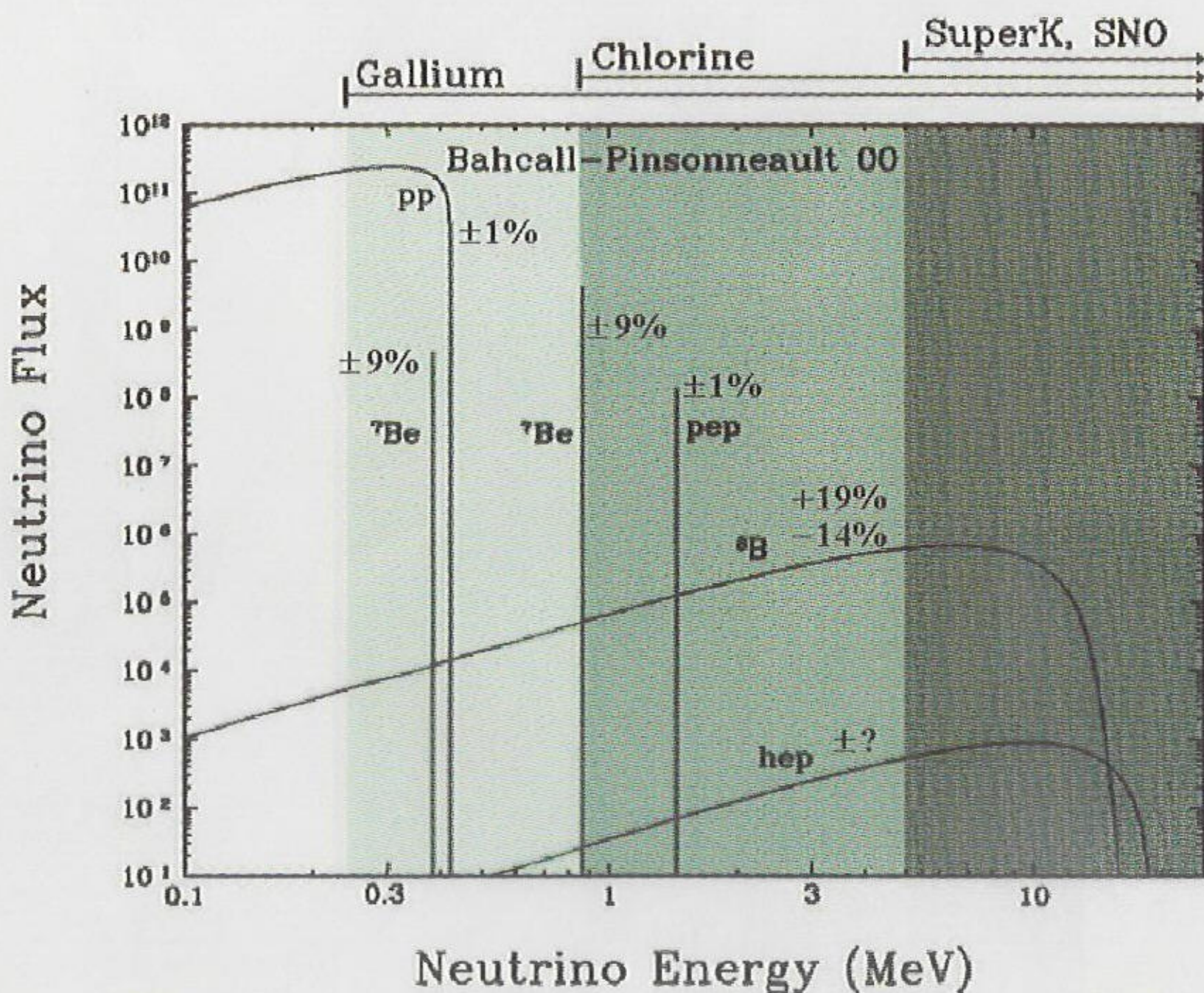
[superallowed transition, simple atom, small energy release $Q = 18.6 \text{ keV}, \dots$]

The spectrum is determined by phase space

$$\frac{dN}{dT} = C p E (Q-T) \sqrt{(Q-T)^2 - m_\nu^2} F(E)$$

m_ν is mass of ν_1 Fermi function
["kinks" at lower E 's for $\nu_{2,3}$]

Experimental spectrum is fitted by m_ν^2 and many other parameters (Q , background term, normalization, \dots) \Rightarrow From Present Day Experiments, **NO INDICATION IN FAVOUR OF $m_\nu \neq 0$**



- To observe ν -oscillations,

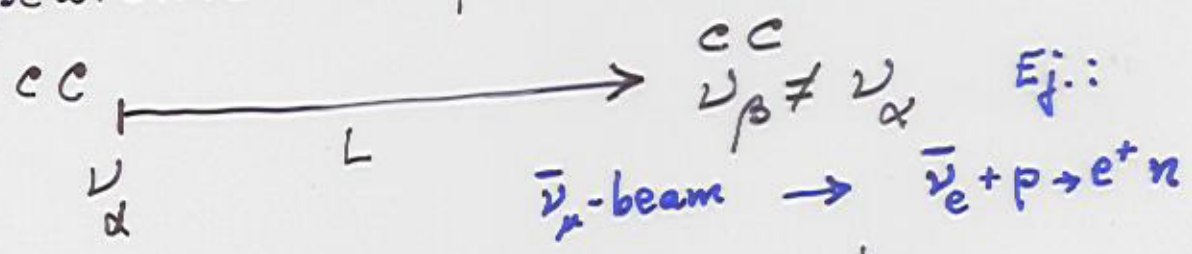
1) At least one Δm^2

$$\Delta m^2 \gtrsim \frac{P}{L}$$

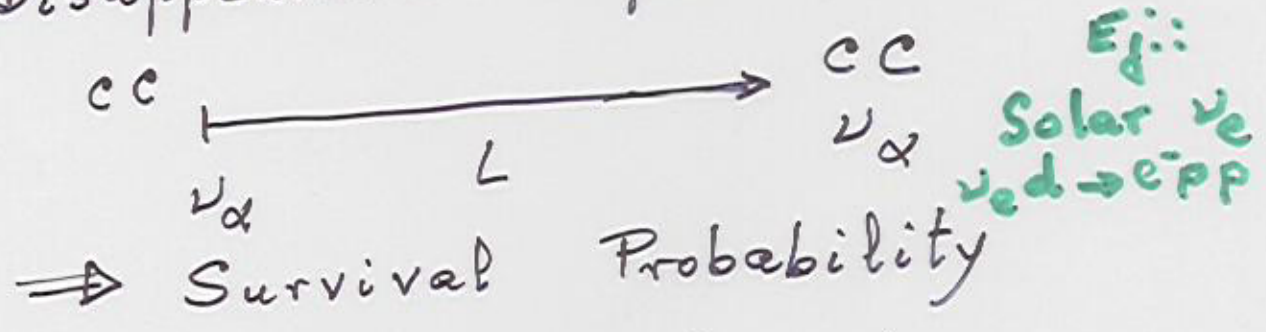
2) Mixing of Prepared and Detected Flavours to common Mass-eigenstates

- Types of Experiments:

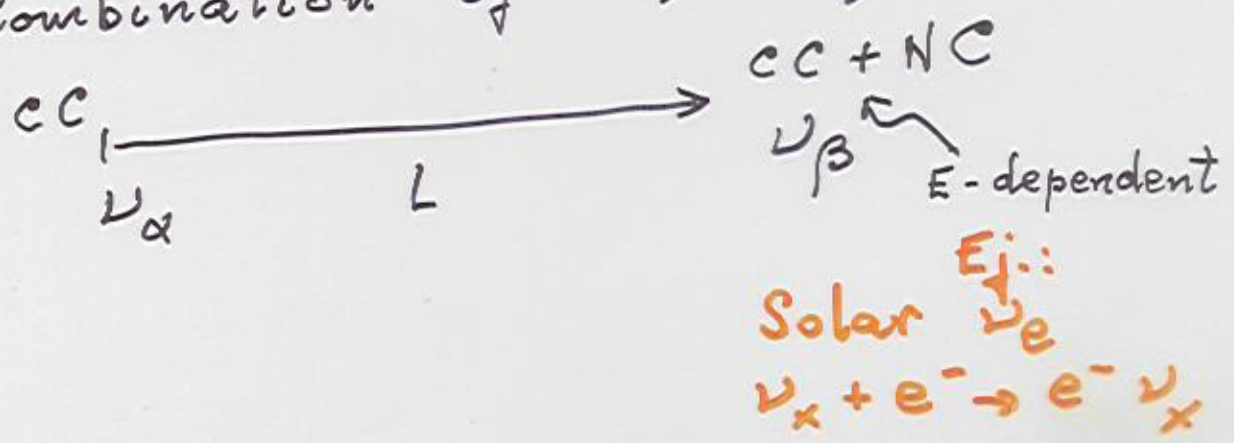
1) Appearance Experiments



2) Disappearance Experiments



3) Combination of 1) + 2)



- SYMMETRIES

$$\nu \rightarrow \bar{\nu} \Rightarrow U \rightarrow U^*$$

$$CPT \rightsquigarrow A(\bar{\alpha} \rightarrow \bar{\beta}; t) = A^*(\alpha \rightarrow \beta; -t)$$

$$T \rightsquigarrow |A(\beta \rightarrow \alpha; t)| = |A(\alpha \rightarrow \beta; t)|$$

$$|A(\bar{\beta} \rightarrow \bar{\alpha}; t)| = |A(\bar{\alpha} \rightarrow \bar{\beta}; t)|$$

$$CP \rightsquigarrow |A(\bar{\alpha} \rightarrow \bar{\beta}; t)| = |A(\alpha \rightarrow \beta; t)|$$

i) ~~CP~~, ~~T~~ in Appearance only

$$A^*(\alpha \rightarrow \alpha; -t) = A(\alpha \rightarrow \alpha; t)$$

In $\alpha \rightarrow \alpha$, $CPT \Rightarrow CP$ Equality

ii) CP-odd Probability

$$D_{\alpha\beta} \equiv |A(\alpha \rightarrow \beta; t)|^2 - |A(\bar{\alpha} \rightarrow \bar{\beta}; t)|^2$$

UNIQUE for 3 Flavours

$$\boxed{D_{e\mu} = D_{\mu\tau} = D_{\tau e}}$$

iii) T-odd Probability

$$T_{\alpha\beta} \equiv |A(\alpha \rightarrow \beta; t)|^2 - |A(\beta \rightarrow \alpha; t)|^2$$

$$\xrightarrow{CPT} |A(\alpha \rightarrow \beta; -t)|^2$$

is ODD function of t (L).

The Pontecorvo MNS Matrix

$$\begin{Bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{Bmatrix} = U \begin{Bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{Bmatrix} \quad \Leftrightarrow$$

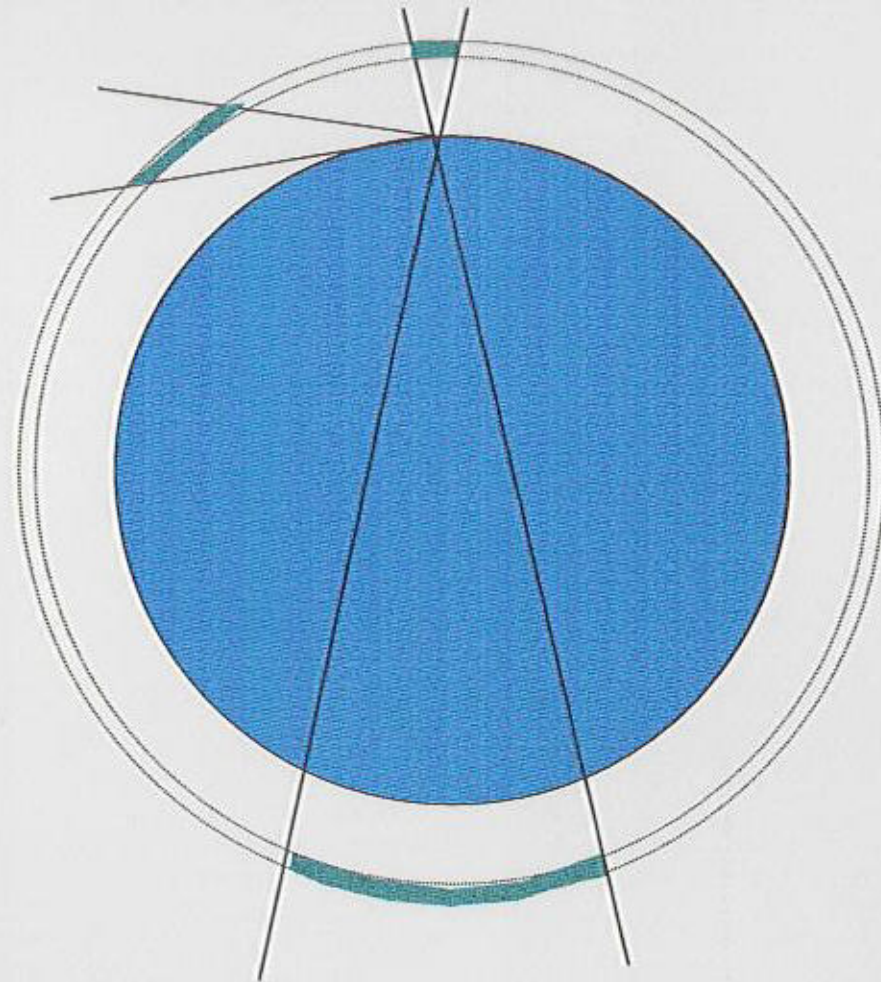
For Flavour oscillations
 U : 3 mixings, 1 phase

$$U = \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}}_{\text{Atmospheric LBL-beams}} \underbrace{\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}}_{\text{Atmospheric (through Matter)? Reactor? Appearance } \nu_\mu \rightarrow \nu_e!} \underbrace{\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{\text{Solar Reactor (if LMA-MSW)}}$$

Atmospheric
 LBL - beams

? ?
 Atmospheric
 (through Matter)?
 Reactor ?
 Appearance $\nu_\mu \rightarrow \nu_e!$

Solar
 Reactor (if LMA-MSW)



Up-Down Symmetry **does not imply** isotropy:

$$V_{\text{v source}}(\theta+\delta\theta) \propto l^2(\theta)/\cos\theta_{\text{emission}}$$

need for 2 necessary ingredients:

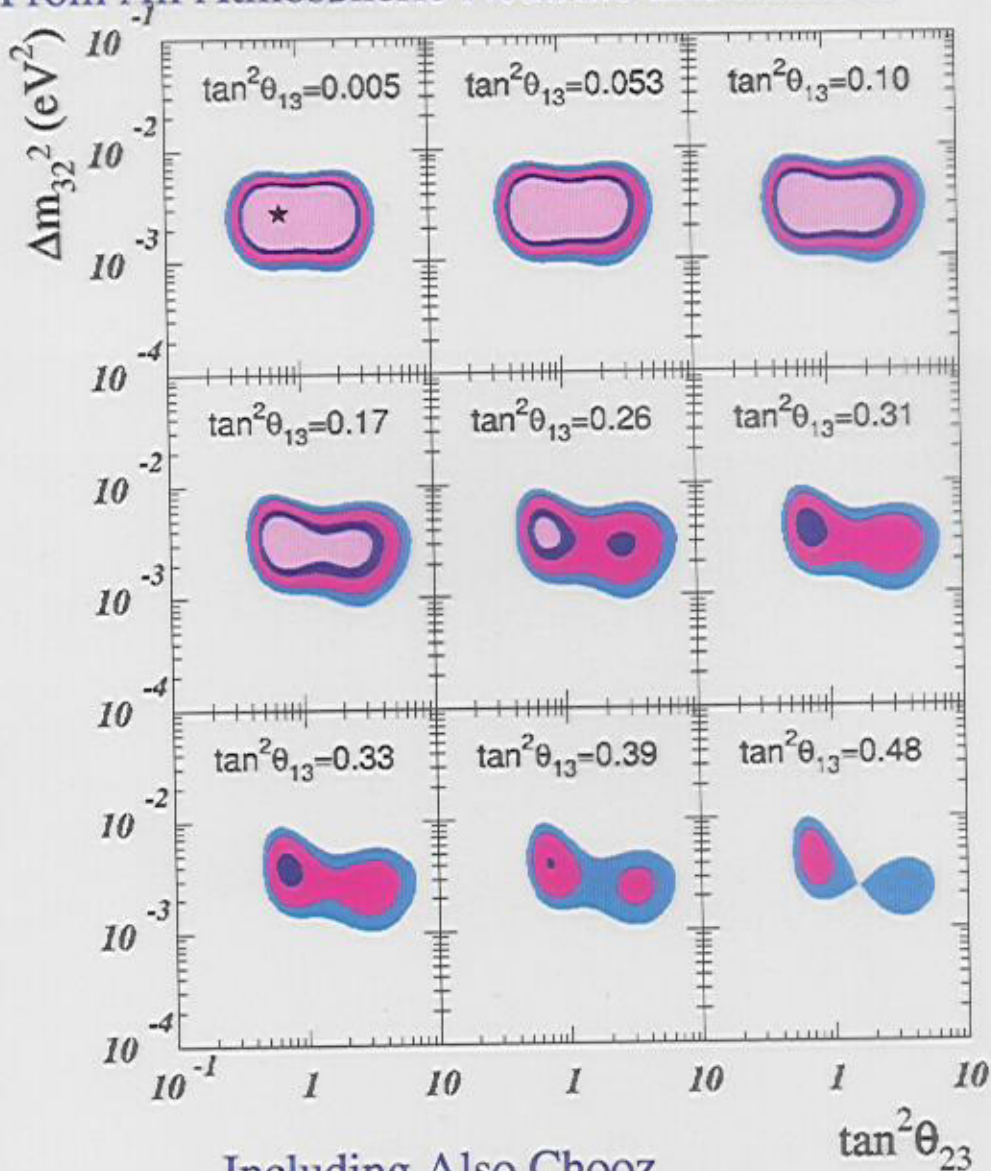
- 1) 3D calc with the whole spherical geometry of earth
- 2) consider the $\langle\theta_{\text{emission}}^2\rangle \neq 0$ [after particle production, decay and transport]

➡ **Enhancement of flux on the horizontal direction**

3- ν Atmospheric Neutrino Oscillation Parameters

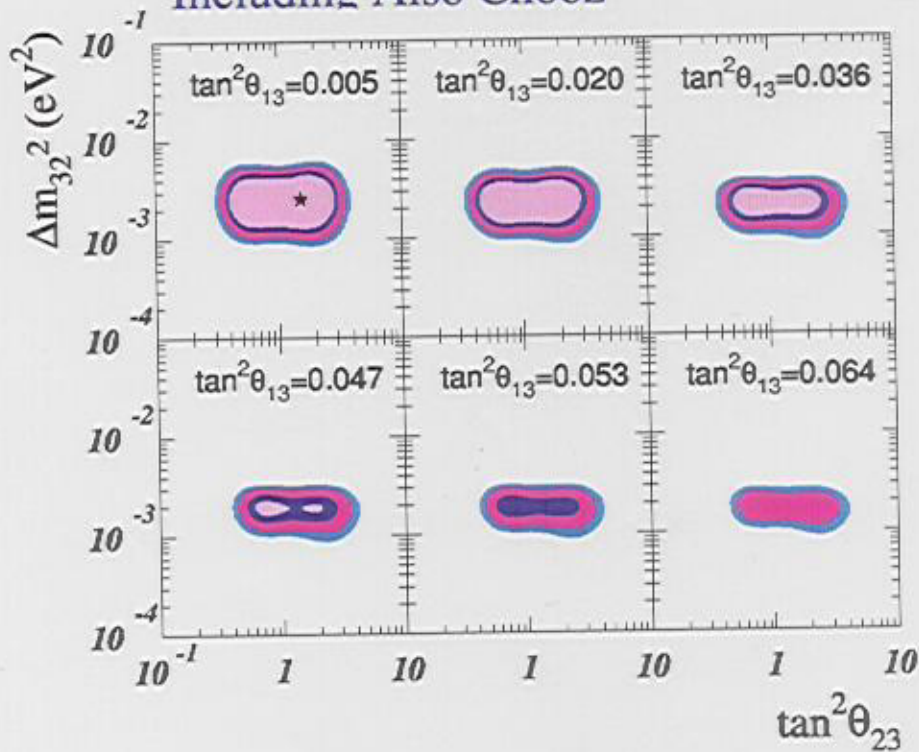
Update of M.C.G-G, M. Maltoni, C. Peña-Garay, J. Valle, PRD63 (2001)

From All Atmospheric Neutrino Experiments



Concha González-García

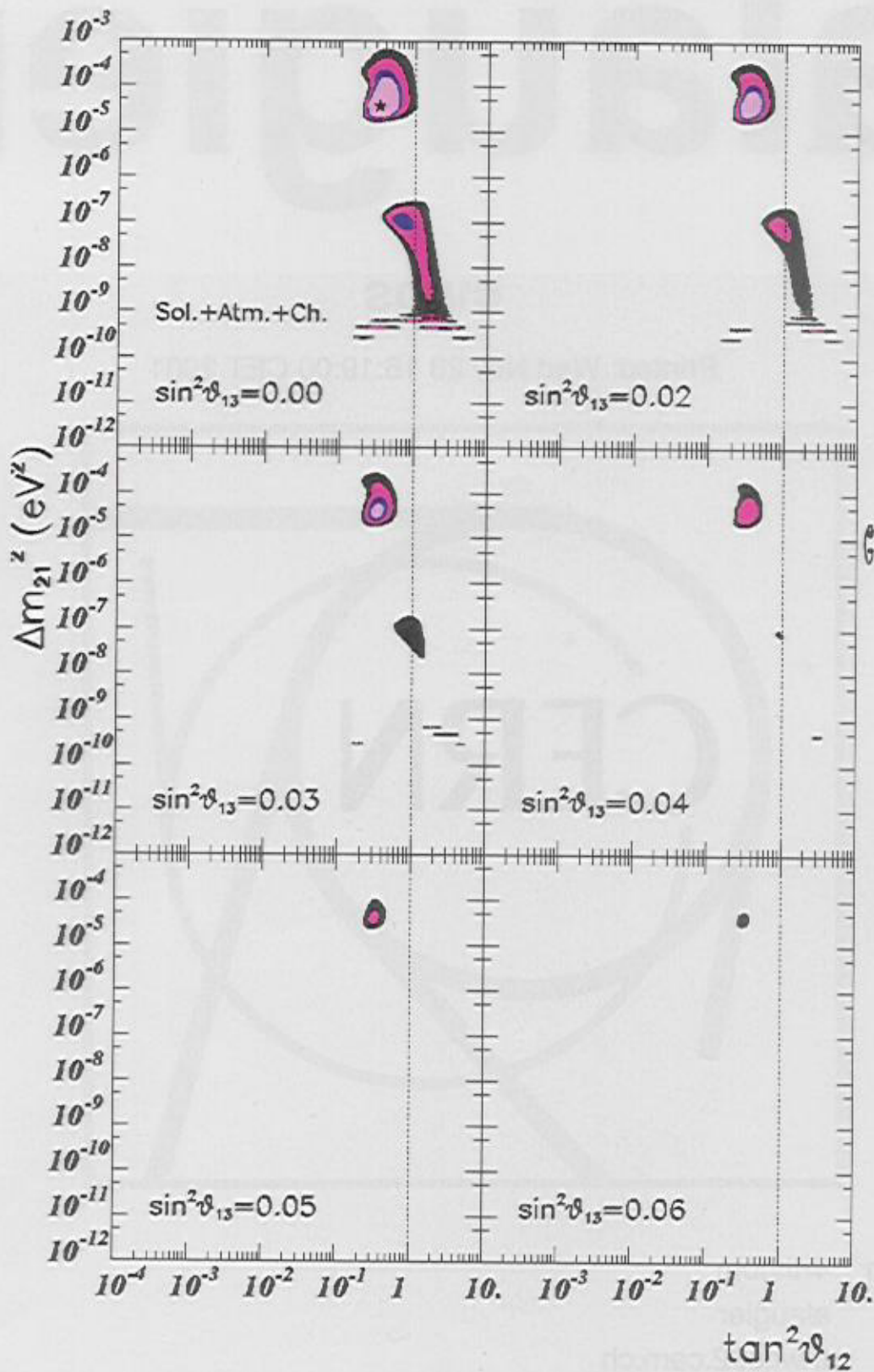
Including Also Chooz



Concha

3- ν Combined Analysis

Effect of Reactor + Atmospheric Data in Solar Regions

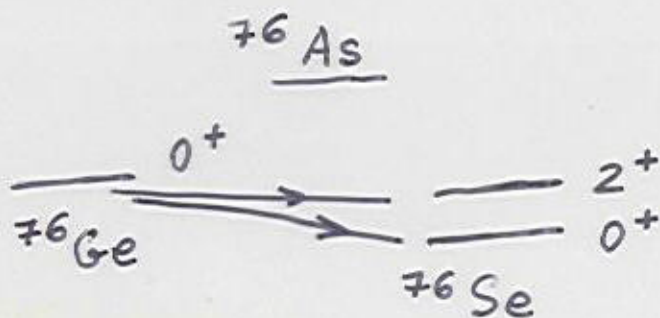
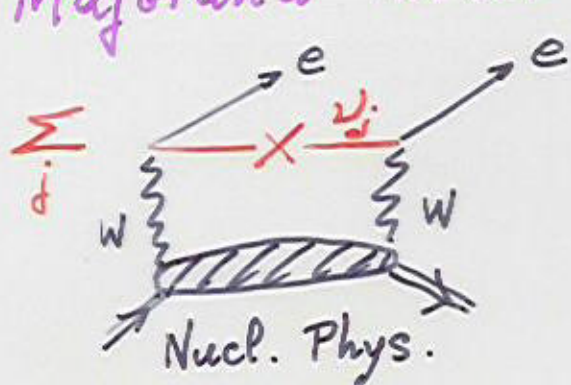


The best way to test Dirac or Majorana is to search for $\beta\beta_{0\nu}$:

$$(A, Z) \rightarrow (A, Z+2) + e^- + e^-,$$

allowed if L is not conserved.

Majorana mass term $\Rightarrow \Delta L = 2$.



The propagator is here

$$\overbrace{\nu_{eL}(x_1) \nu_{eL}^T(x_2)} =$$

$$= - \sum_i U_{ei}^2 \frac{1-\gamma_5}{2} \chi_i(x_1) \bar{\chi}_i(x_2) \frac{1-\gamma_5}{2} C$$

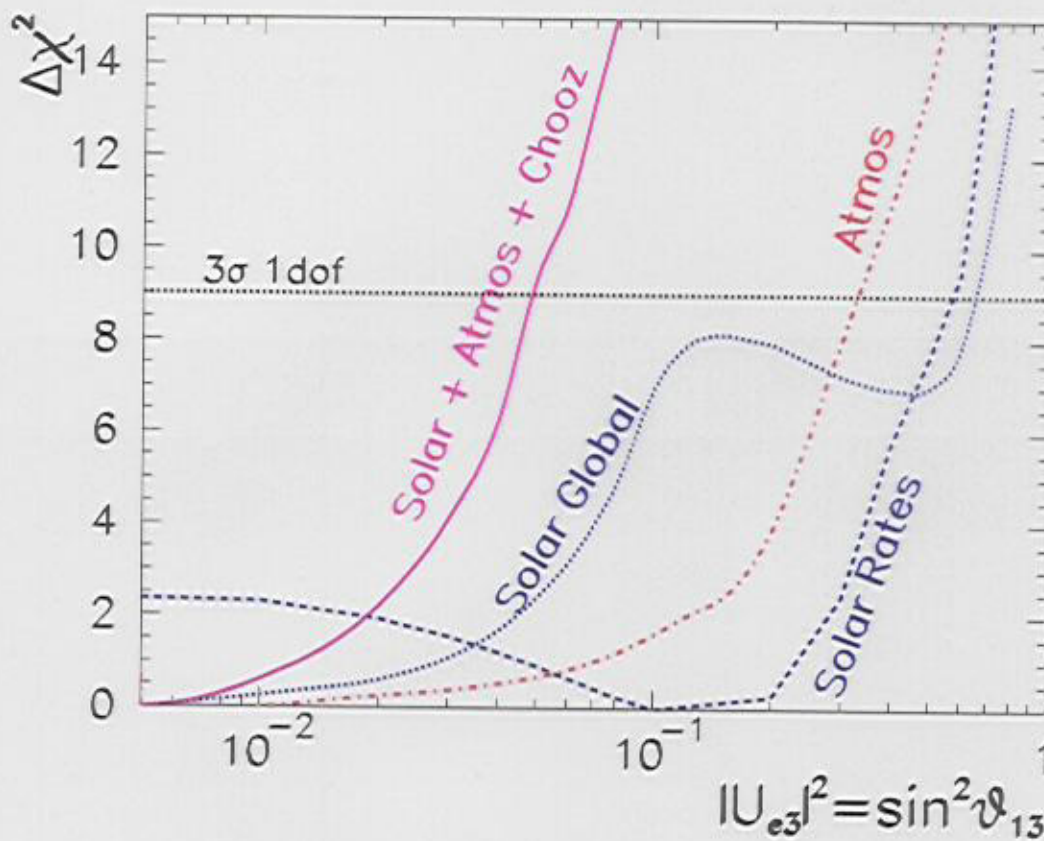
$$= \sum_i U_{ei}^2 m_i \frac{(-i)}{(2\pi)^4} \int d^4p \frac{e^{ip(x_1-x_2)}}{p^2 + m_i^2} \frac{1-\gamma_5}{2} C$$

If m_i are small ($< \text{MeV}$), they can be neglected in the propagator

$$\Rightarrow \text{Amp}[\beta\beta_{0\nu}] = \langle m_\nu \rangle (\text{Ph. Sp.})(\text{Nucl. Phys.})$$

$$\langle m_\nu \rangle = \sum_i U_{ei}^2 m_i$$

3-ν Combined Analysis: Constraints on θ_{13}



Concha
González-García

• 3σ (1dof) allowed ranges of parameters

$$\begin{aligned}
 1.4 \times 10^{-3} < \Delta m_{32}^2 / \text{eV}^2 < 6 \times 10^{-3} \\
 1.9 \times 10^{-5} < \Delta m_{21}^2 / \text{eV}^2 < 2.7 \times 10^{-4} \\
 0.4 < \tan^2 \theta_{23} < 3. \\
 0.22 < \tan^2 \theta_{12} < 0.71 \\
 \sin^2 \theta_{13} < 0.05
 \end{aligned}$$

Mixing Matrix U

$$U = \begin{pmatrix} 0.74 - 0.91 & 0.41 - 0.65 & < 0.34 \\ 0.21 - 0.64 & 0.25 - 0.77 & 0.51 - 0.87 \\ 0.05 - 0.56 & 0.40 - 0.83 & 0.48 - 0.85 \end{pmatrix}$$

- KNOWN FACTS (with care)

→ Sterile ν 's not needed in the Mixing

→ $|\Delta m^2_{12}| \ll |\Delta m^2_{23}|$
Solar Atmospheric

→ Atmospheric ν oscillation compatible with $\nu_\mu \rightarrow \nu_\tau$: near maximal mixing

$$c_{23} \sim s_{23} \sim \frac{1}{\sqrt{2}}$$

$$|\Delta m^2_{23}| \sim (1.5 - 4) \times 10^{-3} \text{ eV}^2$$

→ Solar ν oscillation compatible with LMA-MSW: near maximal mixing $\nu_e \rightarrow \nu'$

$$c_{12} \sim s_{12} \sim \frac{1}{\sqrt{2}}$$

$$|\Delta m^2_{12}| \sim (2 - 30) \times 10^{-5} \text{ eV}^2$$

→ CHOOZ reactor + Atmospheric

$$s^2_{13} \lesssim 0.05$$

- For $m_1 \ll m_2 \ll m_3$, $\Delta m_{12}^2 \frac{R}{2P} \ll 1$, 12

$$P_{\alpha \rightarrow \alpha'} = 2 |\tilde{U}_{\alpha'3}|^2 |\tilde{U}_{\alpha 3}|^2 \left(1 - \cos \Delta m_{31}^2 \frac{R}{2P}\right)$$

$\alpha' \neq \alpha$

\Rightarrow **SAME FORM** as for 2ν , BUT
different meaning of "mixing"

$$P_{\alpha\alpha} = 1 - \sum_{\beta \neq \alpha} P_{\alpha\beta} =$$

$$= 1 - 2 |\tilde{U}_{\alpha 3}|^2 (1 - |\tilde{U}_{\alpha 3}|^2) \left(1 - \cos \Delta m_{31}^2 \frac{R}{2P}\right)$$

\Rightarrow Disappearance Reactor Experiment

gives $|\tilde{U}_{e3}|^2$: CHOOZ
Palo Verde

→ Zeroth Order Approximation

$$U = \begin{pmatrix} \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} & 0 \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{2} & \frac{1}{2} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

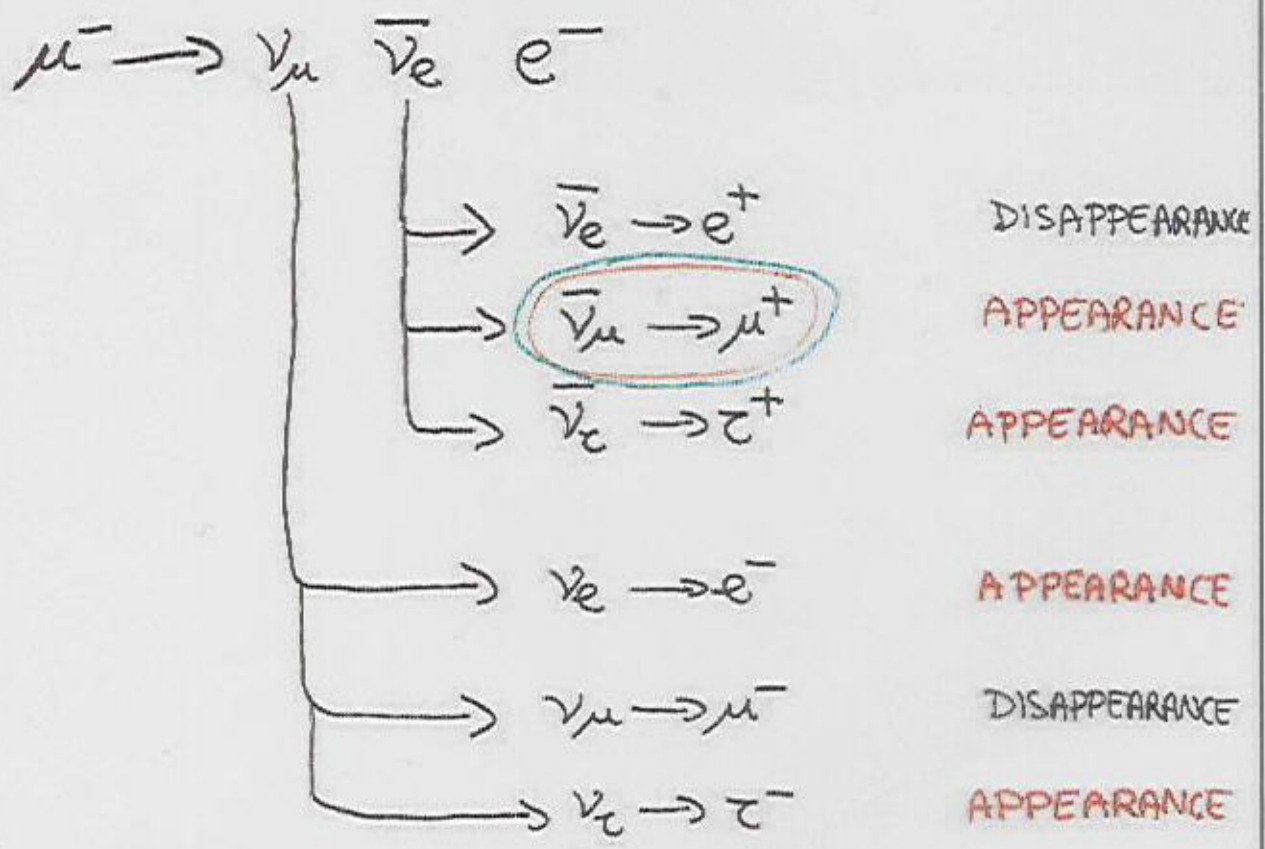
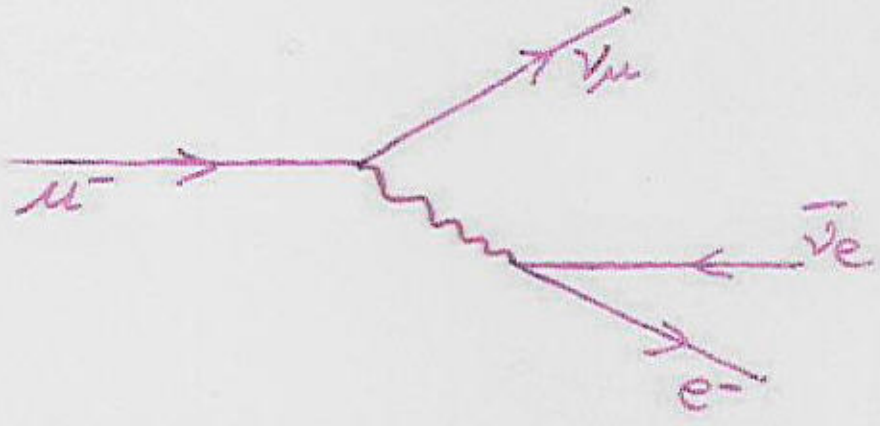
Main Question: V_{e3} ?

[to be answered before ν Factories]

⇒ No ~~CP~~ - physics without $V_{e3} \neq 0!$

ν' FACTORIES

BEAM COMPOSITION



SEARCH FOR WRONG-SIGN μ

EXTREMELY SENSITIVE APPEARANCE MEASUREMENT
(Geer)

UNKNOWNNS

- Dirac or Majorana?

$m_\nu \neq 0$ without ν_R ? $\rightarrow M_j$

Dimension 5 - Effective \mathcal{L}
in $SU(2) \times U(1)$

"See-Saw" Mechanism

$$\beta\beta_{0\nu} \rightarrow \langle m_\nu \rangle \equiv \sum_i U_{ei}^2 m_i$$

- Absolute Masses

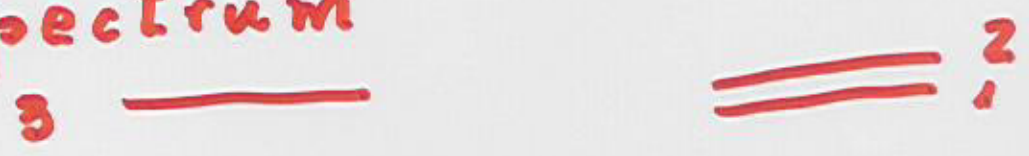
Δm_{jk}^2 \oplus The most sensitive
"Direct" Mass measure

\rightsquigarrow ${}^3\text{H}$ Beta Decay and/or SN

$$m_\beta^2 \equiv \sum_i |U_{ei}|^2 m_i^2$$

\rightsquigarrow $\langle m_\nu \rangle$ in $\beta\beta_{0\nu}$

- Spectrum



vs.



- Follow the fathers of SM :

(AT PRESENT ENERGIES) $SU(2) \times U(1)$

with a particle content WITHOUT ν_R

- Is it possible (with only ν_L) to have $m_\nu \neq 0$?

A priori, YES \longleftrightarrow Majorana

$$\mathcal{L}^{\text{Maj}}_{\text{mass}} = -\frac{1}{2} \bar{\nu}_L M \nu_L^c + \text{h.c.}$$

$SU(3)_{\text{colour}} \otimes U(1)_{\text{e.m.}}$ invariant \Rightarrow legal

$$\mathcal{L}^{\text{Maj}} + \text{Pauli} \Rightarrow M^T = M$$

Diagonalization \rightarrow

$$M = U m U^T$$

$$\mathcal{L}^{\text{Maj}}_{\text{mass}} = -\frac{1}{2} \bar{X} m X$$

$$X \equiv \nu^+ \nu_L + (\nu^+ \nu_L)^c$$

Majorana Field

$$X = X^c = C \bar{X}^T$$

\Rightarrow Neutrinos would be "true" neutrals

- NO conserved GLOBAL LEPTON NUMBER

$\Rightarrow \beta\beta_{\nu\nu}$ allowed

BUT ...

- $\int \mathcal{L}_{\text{mass}}^{\text{Maj}}$ cannot be GENERATED by
SSB of a RENORMALIZABLE (dim. 4)
Yukawa interaction! Higgs triplet?

1970's \rightarrow 1990's: Requirement of
RENORMALIZABILITY has a different
PHILOSOPHY:

- Take particle content of SM and ask
WHAT IS THE LOWEST DIMENSION (NON-
RENORMALIZABLE) OPERATOR WITH
 $SU(2) \times U(1)$ GAUGE INVARIANCE?

UNIQUE
dim. 5



$$\mathcal{L}_{\text{eff}} = -\frac{1}{2\Lambda} (\bar{\tilde{l}}_L \psi) F (\tilde{\psi}^+ l_L)$$

$$\tilde{l} = i \tau_2 l^c = i \tau_2 C \bar{l}^T$$

- The FIRST window to BEYOND SM

- $F = F^T$ matrix in flavour-space

- Coupling $\frac{1}{\Lambda} \Leftrightarrow \Lambda \equiv$ New Physics Scale

\Rightarrow 1) $\Delta L = 2$

2) $F \Rightarrow$ Mixing, LFV

3) After **SSB** : $\mathcal{L}_{eff} \rightarrow \mathcal{L}^{Maj\ mass}$
 $M = \frac{v^2}{\Lambda} F$ "see-saw" type

Conclusion

$\mathcal{L}^{Maj\ mass}$ for "light" neutrinos **GENERATED** from \mathcal{L}_{eff} (at present energies).

- Origin of Λ ?

1) Heavy mass ν_R ?

2) Fierz - reordering

$$\mathcal{L}_{eff} = -\frac{1}{4\Lambda} (\bar{l}_L F \sum l_L) (\bar{\psi} + \sum \psi)$$

Higgs triplet ?

- A particular example : **SO(10) GUT**

\Rightarrow As part of **SB**, ν_R obtains a large Majorana mass Λ .

- The Dirac mass terms $\nu_R \leftrightarrow \nu_L m_D$ produce mixing of heavy ν_R into $\nu_L \Rightarrow m_\nu = \frac{m_D^2}{\Lambda}$

"SEE-SAW" MECHANISM

$SU(2) \times U(1)$ gauge invariance
 $SM + \nu_R$ in particle content

\Rightarrow Not only L-R Dirac Mass,
 R Majorana Mass too

$$\begin{aligned}
 \mathcal{L}^{D-M} &= -m_D \bar{\nu}_R \nu_L - \frac{1}{2} m_R \bar{\nu}_R (\nu_R)^c + h.c. \\
 &= -\frac{1}{2} (\bar{n}_L)^c M n_L + h.c.
 \end{aligned}$$

$$n_L \equiv \begin{pmatrix} \nu_L \\ (\nu_R)^c \end{pmatrix}, \quad M = \begin{pmatrix} 0 & m_D \\ m_D & m_R \end{pmatrix}$$

\rightarrow For one family and $m_D \ll m_R$,

$$m_1 \approx \frac{m_D^2}{m_R}, \quad m_2 \approx m_R,$$

$$\theta \approx \frac{m_D}{m_R}; \quad \eta_1 = -1, \eta_2 = 1$$

$$\Rightarrow \boxed{\nu_L \approx -i \nu_{1L}, \quad (\nu_R)^c \approx \nu_{2L}}$$

ν_1, ν_2 Majorana Fields

- Majorana Fields

$$\nu_1 \simeq i\nu_L - i(\nu_L)^c; \quad \nu_2 = \nu_R + (\nu_R)^c$$

1 light neutrino $m_1 \ll m_D$

+
1 heavy neutrino $m_2 \gg m_D$

$\rightarrow m_L = 0 \iff$ Global L violated
only by $-\frac{1}{2} m_R \bar{\nu}_R (\nu_R)^c$
characterized by large mass.

$\rightarrow m_D \simeq \mathcal{O}(\text{quark, lepton}) \text{ mass}$

SMALLNESS OF NEUTRINO MASS

\iff L-VIOLATION at $m_R \simeq M_{\text{GUT}}$

- For 3 families, $m_i \simeq \frac{(m_f^i)^2}{M_i} \ll m_f^i$

- 1) ν_α with definite mass \rightarrow Majoranas
- 2) 3 "active" massive neutrinos
- 3) Hierarchy of masses $m_1 \ll m_2 \ll m_3$.

$$H = \frac{1}{2E} \left\{ U \begin{pmatrix} 0 & & \\ & \Delta m_{12}^2 & \\ & & \Delta m_{13}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} a & & \\ & 0 & \\ & & 0 \end{pmatrix} \right\}$$

→ Interference with MEDIUM Effects sees the SIGN.

→ $\nu \rightarrow \bar{\nu} : U \rightarrow U^*, a \rightarrow -a$

Suggestion:

Muon-Charge Asymmetry for Atmospheric Neutrinos from below

$$\frac{(\nu_{atm} \rightarrow \nu_\mu) - (\bar{\nu}_{atm} \rightarrow \bar{\nu}_\mu)}{+}$$

→ Magnetized Iron: MINOS, MONOLITH, ...

But ... ≈ 100 Kton

→ U_{e3} :

Appearance $\nu_e \rightarrow \nu_\mu \sim U_{e3}$
better than Disappearance

→ δ : ν FACTORIES

$$(\nu_e \rightarrow \nu_\mu) - (\bar{\nu}_e \rightarrow \bar{\nu}_\mu) \leftrightarrow \cancel{a} \oplus a \neq 0$$

or SUPERBEAMS ($\nu_\mu \rightarrow \nu_e$)

CONCLUSIONS

- Exciting times in ν ' Physics.
- Convincing and Consistent Evidence for ν ' Oscillation in

$$\text{Atmospheric : } \Delta m_{23}^2 \oplus \theta_{23} \approx \frac{\pi}{4}$$

$$\text{Solar : } \Delta m_{12}^2 \oplus \theta_{12} \approx \frac{\pi}{4}$$

- Test of Atmospheric by LBL.

- Test of Solar by BOREXINO plus Reactor Exp. KAMLAND.

- $\theta_{13} \rightarrow$ Appearance $\nu_{\mu} \leftrightarrow \nu_e$.

- $\delta \rightarrow$ ~~CP~~, \uparrow

- Sign $[\Delta m^2] \leftrightarrow$ Interference with Medium Effects

- Above all,

THEORY expects ν 's are MAJORANA particles